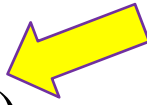


Fuel Depletion

$$N \sim 10^{22} \text{ cm}^{-3}, \phi \sim 10^{14} \text{ cm}^{-2} \text{ s}^{-1}$$

**Time scale:
Days and months.**

- More depletion ► increase steady state flux by means of reducing absorbers.

- For a given fuel isotope $\frac{\partial N_f(\vec{r}, t)}{\partial t} = -N_f(\vec{r}, t)\sigma_a^f\phi(\vec{r}, t)$ 

- For **constant flux** ϕ_0 the solution is **Exponential burnup**

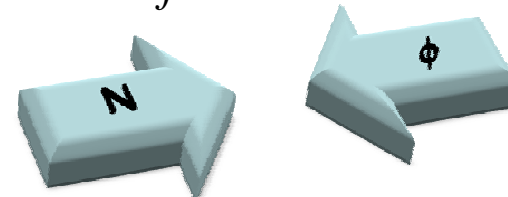
$$N_f(\vec{r}, t) = N_f(\vec{r}, 0)e^{-\sigma_a^f \phi_0(\vec{r})t} = N_f(\vec{r}, 0)e^{-\sigma_a^f \Phi(\vec{r}, t)}$$

- For **time varying flux**

$$N_f(\vec{r}, t) = N_f(\vec{r}, 0)e^{-\sigma_a^f \int_0^t \phi(\vec{r}, t') dt'} = N_f(\vec{r}, 0)e^{-\sigma_a^f \Phi(\vec{r}, t)}$$

Neutron fluence

Solve numerically.



Fuel Depletion

- **Constant power.**

$$P(\vec{r}, t) = \overset{\downarrow}{w} N_f(\vec{r}, t) \overset{\uparrow}{\sigma_f^f} \phi(\vec{r}, t) = P(\vec{r}, 0) = P_0(\vec{r})$$

Energy released per fission

 Fission rate

$$N_f(\vec{r}, t)\phi(\vec{r}, t) = N_f(\vec{r}, 0)\phi(\vec{r}, 0)$$

$$\sum_f(\vec{r}, t)\phi(\vec{r}, t) = \sum_f(\vec{r}, 0)\phi(\vec{r}, 0)$$

- Power \sim flux only over short time periods during which N_f is constant.

$$\frac{\partial N_f(\vec{r}, t)}{\partial t} = -N_f(\vec{r}, t) \sigma_a^f \phi(\vec{r}, t) \approx -\frac{P_0(\vec{r})}{w}$$

Linear depletion!

- The solution is obviously

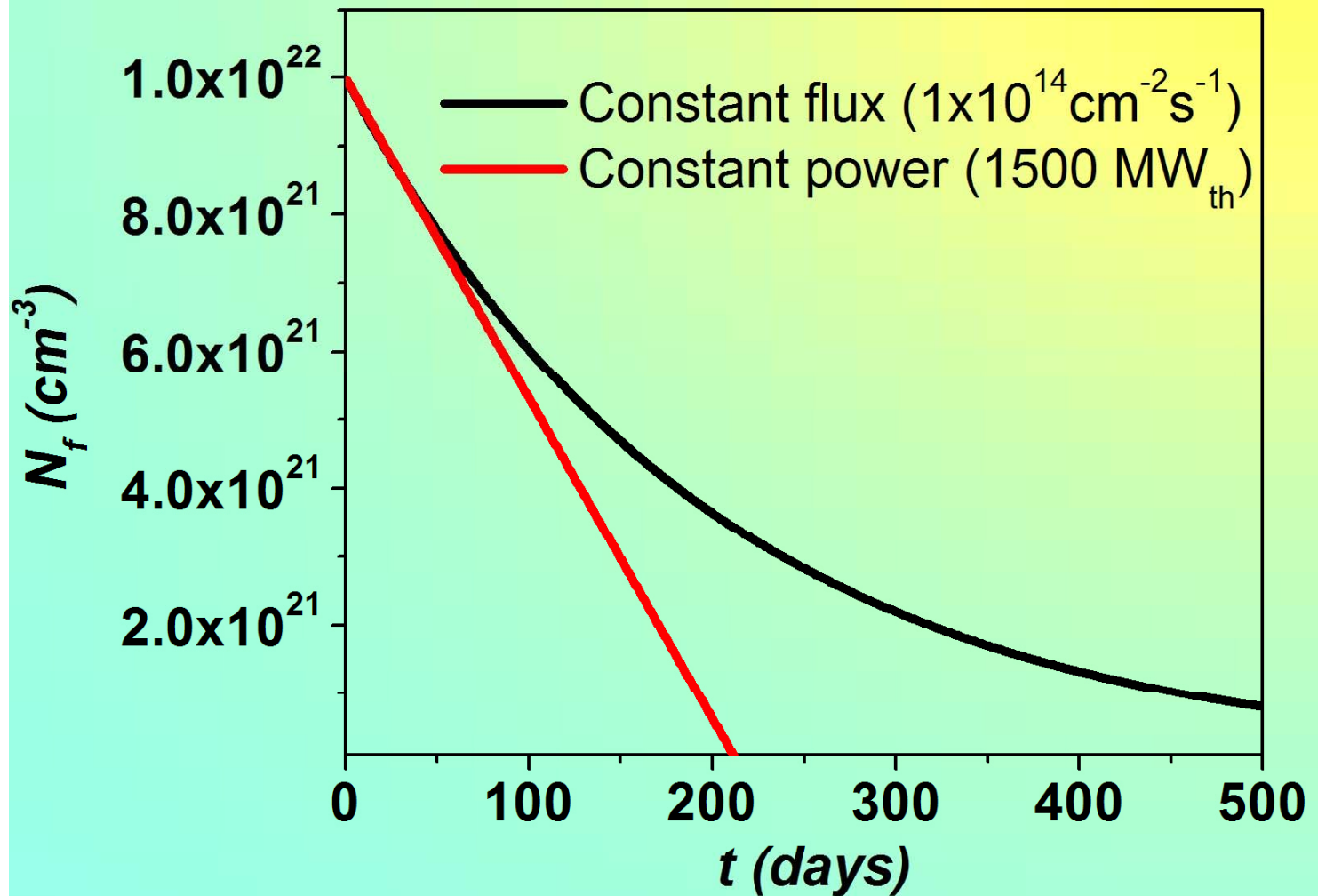
$$N_f(\vec{r}, t) \approx N_f(\vec{r}, 0) - \frac{P_0(\vec{r})}{w} t$$

$$\sigma_f^f \approx \sigma_a^f$$

Fuel Depletion

HW 31

Do the calculations for different flux and power levels.



Poisoning and Fuel Depletion

Infinite, critical homogeneous reactor.

$$k_{\infty} = \epsilon \eta f \rho = \epsilon \eta \rho \frac{\sum_a^f(t)}{\sum_a^f(t) + \sum_a^{clad} + \sum_a^{moderator}(t) + \sum_a^{poison}(t) + \sum_a^{control}(t)}$$

\downarrow \uparrow $\text{thus } \downarrow$

Constant power $N_f(\vec{r}, t) \approx N_f(\vec{r}, 0) - \frac{P_0(\vec{r})}{w} t$

$$\phi(\vec{r}, t) = \frac{N_f(\vec{r}, 0)}{N_f(\vec{r}, t)} \phi(\vec{r}, 0) = \frac{\phi(\vec{r}, 0)}{1 - \sigma_a^f \phi(\vec{r}, 0) t}$$

$$= N_f(\vec{r}, 0) - N_f(\vec{r}, t) \sigma_a^f \phi(\vec{r}, t) t$$

$$= N_f(\vec{r}, 0) - N_f(\vec{r}, 0) \sigma_a^f \phi(\vec{r}, 0) t$$

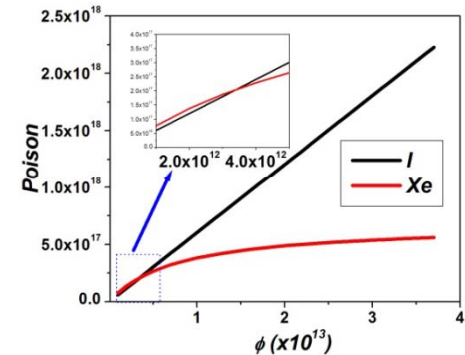
$$= N_f(\vec{r}, 0) [1 - \sigma_a^f \phi(\vec{r}, 0) t]$$

$$\sum_a^f(\vec{r}, t) = \sum_a^f(\vec{r}, 0) [1 - \sigma_a^f \phi(\vec{r}, 0) t]$$

Poisoning and Fuel Depletion

$$Xe(t) = \frac{Xe_\infty}{(\gamma_I + \gamma_{Xe}) \sum_f \phi_0} (1 - e^{-(\lambda_{Xe} + \sigma_a^{Xe} \phi_0)t}) + \frac{\gamma_I \sum_f \phi_0}{\lambda_{Xe} - \lambda_I + \sigma_a^{Xe} \phi_0} (e^{-(\lambda_{Xe} + \sigma_a^{Xe} \phi_0)t} - e^{-\lambda_I t})$$

Constant



$$\sum_a^{Xe}(\vec{r}, t) = \sigma_a^{Xe} Xe_\infty(\vec{r}, t) = \frac{(\gamma_I + \gamma_{Xe}) \sum_f(\vec{r}, 0) \phi(\vec{r}, 0)}{\frac{\lambda_{Xe}}{\sigma_a^{Xe}} + \phi(\vec{r}, t)}$$

Constant

$$\sum_a^{Sm}(\vec{r}, t) \approx \sigma_a^{Sm} \gamma_{Sm} \sum_f(\vec{r}, 0) \phi(\vec{r}, 0) t$$

$$\phi(\vec{r}, t) = \frac{\phi(\vec{r}, 0)}{1 - \sigma_a^f \phi(\vec{r}, 0) t}$$

- Other fission products (poisons) with less capture cross sections.

Poisoning and Fuel Depletion

- Now we know all macroscopic cross sections.

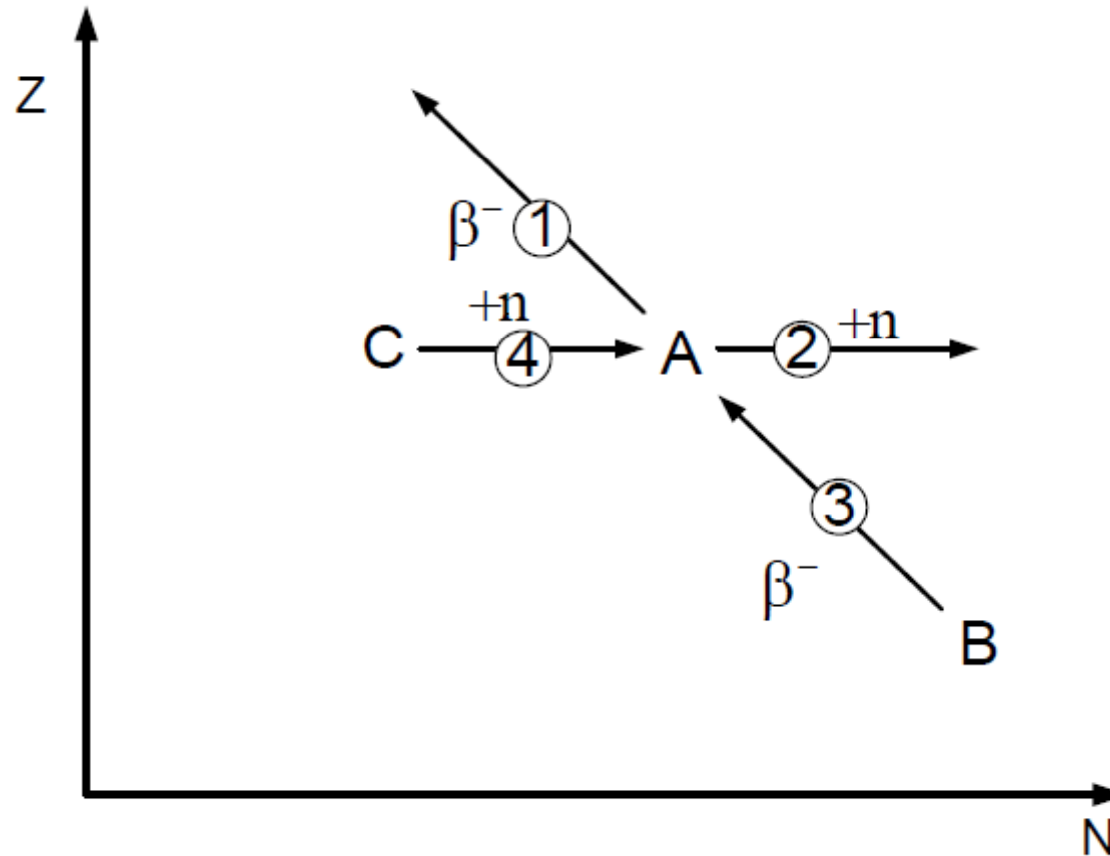
$$k_{\infty} = \varepsilon \eta f \rho = \varepsilon \eta \rho \frac{\sum_a^f(t)}{\sum_a^f(t) + \sum_a^{clad} + \sum_a^{moderator}(t) + \sum_a^{poison}(t) + \sum_a^{control}(t)}$$

↓
↑
↓

- When there are no absorbers left to remove, we need to refuel.
- Absorbers are not only control rods.
- All fuel nuclei should be considered.
- For each species, all sources and sinks should be taken into account.
- Online loading ► environmental.
- ^3H .

Until = 0.
Solve for t to get
upper limit for
“core loading
lifetime”.
 Damaged
 fuel...!

Poisoning and Fuel Depletion



$$\frac{dN_A}{dt} = -\lambda_A N_A - \phi \sigma_A N_A + \lambda_B N_B + \phi \sigma_\gamma^C N_C + F(t)$$

Fuel loading

Poisoning and Fuel Depletion

- **Some poisons are intentionally introduced into the reactor.**

- **Fixed burnable poisons.**

B, Gd.

More uniform distribution than rods, more intentionally localized than shim.

Flatter flux.

- **Soluble poisons (chemical shim) with caution.**

Boric acid (soluble boron, solbor) in coolant.

Boration and dilution.

Scram emergency shutdown (sodium polyborate or gadolinium nitrate).

- **Non-burnable poisons.**

Chain of absorbers or self shielding.

Power shaping.

Delayed Precursors

$$\frac{1}{v_g} \frac{\partial}{\partial t} \phi_g(\vec{r}, t) = \chi_g \sum_{g'=1}^G \nu_{g'} \sum_{fg'}(\vec{r}) \phi_{g'}(\vec{r}, t) + \sum_{g'=1}^G \sum_{sg'g}(\vec{r}) \phi_{g'}(\vec{r}, t) + S_g^{ext}$$

$$- \sum_{ag}(\vec{r}) \phi_g(\vec{r}, t) - \sum_{sg}(\vec{r}) \phi_g(\vec{r}, t) + \vec{\nabla} \cdot D_g(\vec{r}) \vec{\nabla} \phi_g(\vec{r}, t)$$

- For one-group

$$\frac{1}{v} \frac{\partial}{\partial t} \phi(\vec{r}, t) = \nu \sum_f(\vec{r}) \phi(\vec{r}, t) + S^{ext}$$

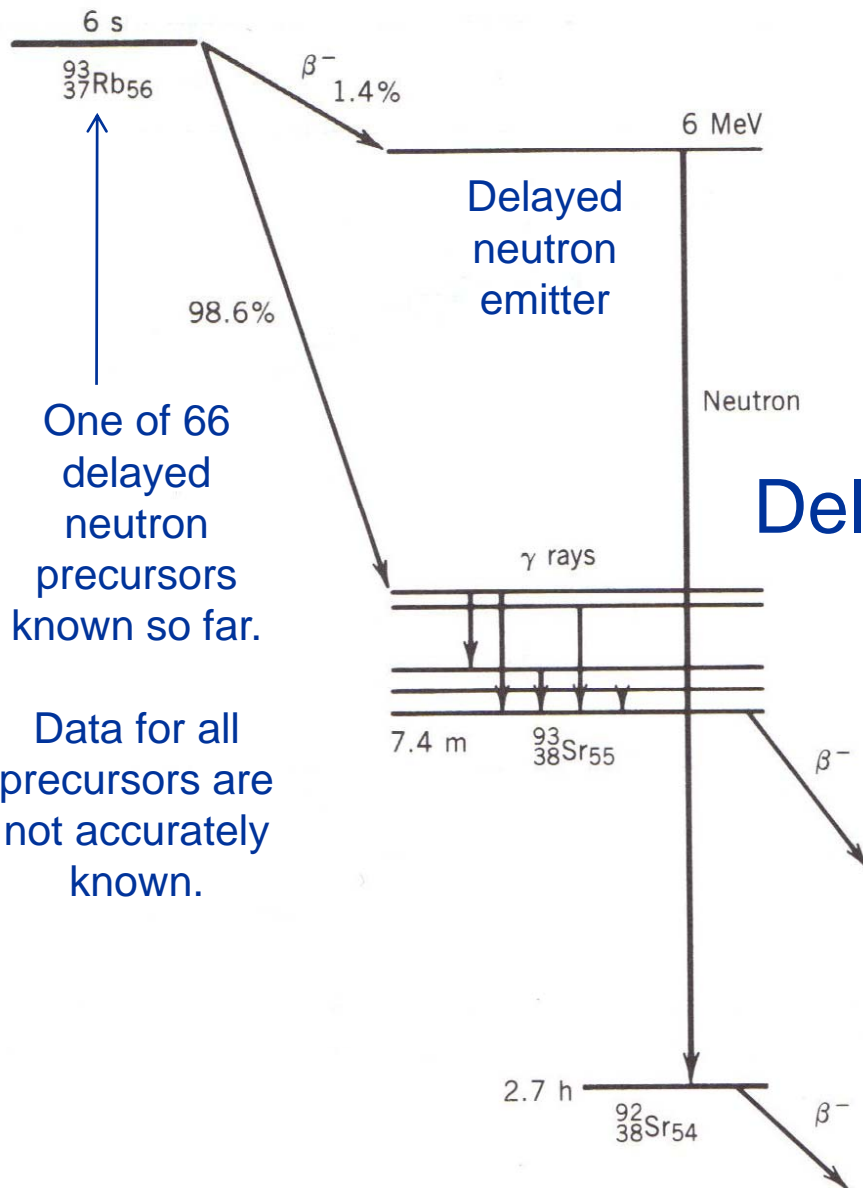
$$- \sum_a(\vec{r}) \phi(\vec{r}, t) + \vec{\nabla} \cdot D(\vec{r}) \vec{\nabla} \phi(\vec{r}, t)$$

- **What about delayed neutrons?**

Delayed Precursors

$$\nu = \nu_p + \nu_d$$

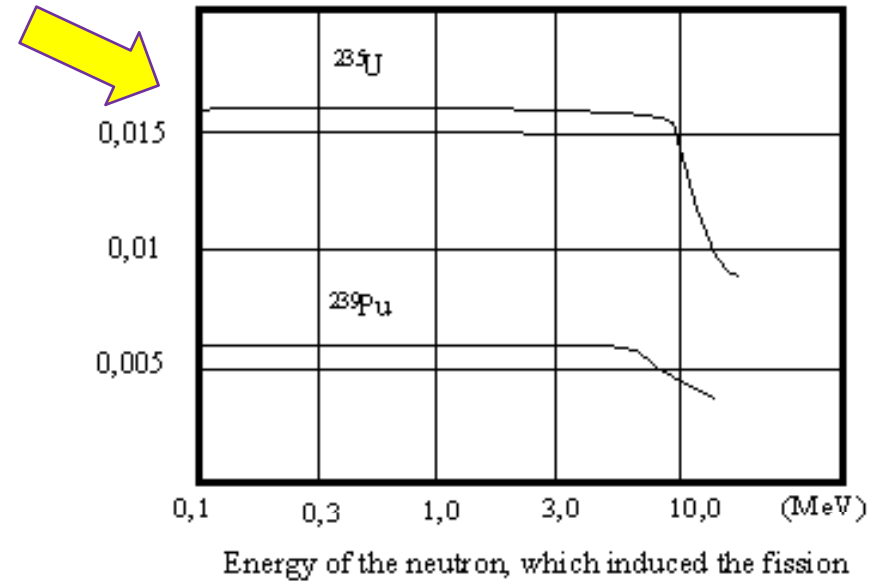
Delayed neutron fraction $\beta = \frac{\nu_d}{\nu}$



One of 66 delayed neutron precursors known so far.

Data for all precursors are not accurately known.

Delayed neutron yield (delayed neutrons / fission)



Delayed Precursors

Fissile nucleus	Delayed neutron / 100 fissions
^{233}U	0.667
^{235}U	1.621
$^{238}\text{U}^*$	4.39
^{239}Pu	0.628
$^{240}\text{Pu}^*$	0.95
^{241}Pu	1.52
$^{242}\text{Pu}^*$	2.21

Increases
with Λ .

**Data for thermal neutron induced fission, except for
* fast neutron induced fission.**

Delayed Precursors

Group #→	1	2	3	4	5	6
$T_{1/2}$ (s)	54.51	21.84	6.0	2.23	0.496	0.179
λ_I	0.0127	0.031	0.1155	0.310	1.397	3.871
β_i/β	.038	0.213	0.188	0.407	0.128	0.026
β_I	0.0002641	0.00148035	0.0013066	0.00282865	0.0008896	0.0001807

^{235}U

ν_p

$\beta < 0.7\% \cong 0.016 / \nu$

$$\frac{1}{\nu} \frac{\partial}{\partial t} \phi(\vec{r}, t) = (1 - \beta) \nu \sum_f(\vec{r}) \phi(\vec{r}, t) + \sum_{i=1}^6 \lambda_i C_i + S^{ext} - \sum_a(\vec{r}) \phi(\vec{r}, t) + \vec{\nabla} \cdot D(\vec{r}) \vec{\nabla} \phi(\vec{r}, t)$$

$$\frac{\partial C_i(\vec{r}, t)}{\partial t} = -\lambda_i C_i(\vec{r}, t) + \beta_i \nu \sum_f(\vec{r}) \phi(\vec{r}, t)$$

Delayed Precursors

- The multi-group equation now becomes

Different energy spectra

$$\frac{1}{v_g} \frac{\partial}{\partial t} \phi_g(\vec{r}, t) = \chi_g^p (1 - \beta) \sum_{g'=1}^G v_{g'} \sum_{fg'}(\vec{r}) \phi_{g'}(\vec{r}, t) + \chi_g^c \sum_{i=1}^6 \lambda_i C_i(\vec{r}, t)$$

$$+ \sum_{g'=1}^G \sum_{sg'}(\vec{r}) \phi_{g'}(\vec{r}, t) + S_g^{ext}$$

$$- \sum_{ag}(\vec{r}) \phi_g(\vec{r}, t) - \sum_{sg}(\vec{r}) \phi_g(\vec{r}, t) + \vec{\nabla} \cdot D_g(\vec{r}) \vec{\nabla} \phi_g(\vec{r}, t)$$

$$\frac{\partial C_i(\vec{r}, t)}{\partial t} = -\lambda_i C_i(\vec{r}, t) + \beta_i \sum_{g'=1}^G v_{g'} \sum_{fg'}(\vec{r}) \phi_{g'}(\vec{r}, t)$$

Full Blown Diffusion Equation

- In steady state

$$\lambda_i C_i(\vec{r}, t) = \beta_i \sum_{g'=1}^G \nu_{g'} \sum_{fg'}(\vec{r}) \phi_{g'}(\vec{r}, t)$$

$$0 = \chi_g^p (1 - \beta) \sum_{g'=1}^G \nu_{g'} \sum_{fg'}(\vec{r}) \phi_{g'}(\vec{r}) + \chi_g^c \beta \sum_{g'=1}^G \nu_{g'} \sum_{fg'}(\vec{r}) \phi_{g'}(\vec{r})$$

$$+ \sum_{g'=1}^G \sum_{sg'}(\vec{r}) \phi_{g'}(\vec{r}) + S_g^{ext} - \sum_{ag}(\vec{r}) \phi_g(\vec{r}) - \sum_{sg}(\vec{r}) \phi_g(\vec{r}) + \vec{\nabla} \cdot D_g(\vec{r}) \vec{\nabla} \phi_g(\vec{r})$$

Significance of χ_g^c depends on whether we have fine or course energy groups.

$$0 = \left[\chi_g^p + (\chi_g^c - \chi_g^p) \beta \right] \sum_{g'=1}^G \nu_{g'} \sum_{fg'}(\vec{r}) \phi_{g'}(\vec{r}) + \sum_{g'=1}^G \sum_{sg'}(\vec{r}) \phi_{g'}(\vec{r}) + S_g^{ext} - \sum_{ag}(\vec{r}) \phi_g(\vec{r}) - \sum_{sg}(\vec{r}) \phi_g(\vec{r}) + \vec{\nabla} \cdot D_g(\vec{r}) \vec{\nabla} \phi_g(\vec{r})$$