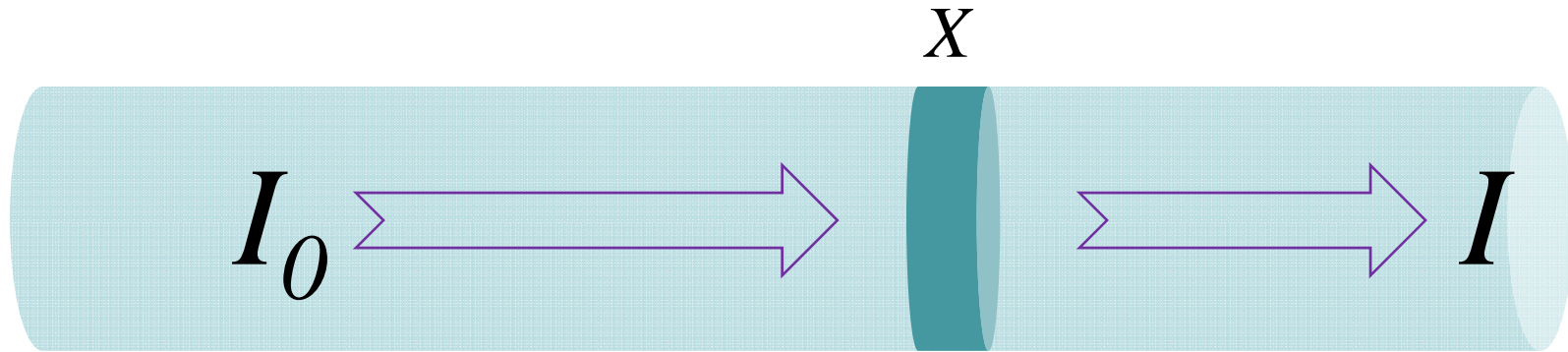


# Neutron Attenuation



Recall  $\Sigma_t = N \sigma_t$

Probability per  
unit path  
length.

$$I(X) = I_0 e^{-\Sigma_t X}$$

mfp for scattering  $\lambda_s = 1/\Sigma_s$

mfp for absorption  $\lambda_a = 1/\Sigma_a$

.....

total mfp  $\lambda_t = 1/\Sigma_t$

**Probability**

$$P_{\text{no-interaction}}(X) = e^{-\Sigma_t X}$$

$$P_{\text{interaction}}(X) = 1 - e^{-\Sigma_t X}$$

# Neutron Moderation (revisited)

Show that, after **elastic** scattering the ratio between the final neutron energy  $E'$  and its initial energy  $E$  is given by:

**HW 6**

$$\frac{E'}{E} = \frac{A^2 + 1 + 2A \cos \theta^{CM}}{(A+1)^2} = \frac{[\cos \theta + \sqrt{A^2 - \sin^2 \theta}]^2}{(A+1)^2}$$

**$^1\text{H}$  ?**

For a head-on collision:  $\left(\frac{E'}{E}\right)_{\min} = \left(\frac{A-1}{A+1}\right)^2 \equiv \alpha$  **Collision Parameter**

After  $n$  **s-wave** collisions:  $\ln E'_n = \ln E - n\zeta$   
 where the average change in **lethargy**  $u = \ln(E_M / E)$

is


$$\overline{\Delta u} = \zeta = \left[ \ln \frac{E}{E'} \right]_{av} = 1 + \frac{(A-1)^2}{2A} \ln \frac{A-1}{A+1}$$

Reference

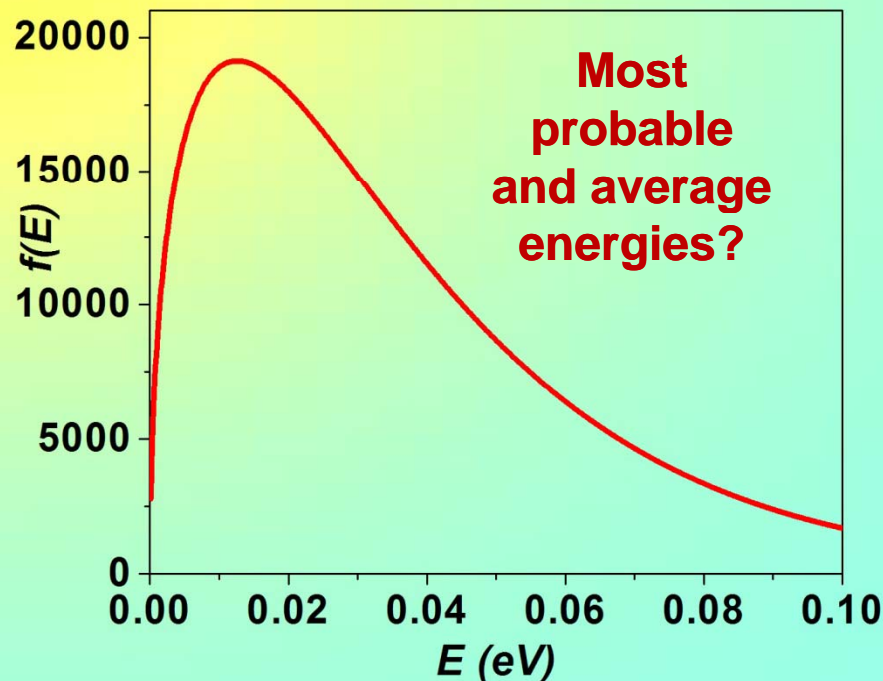
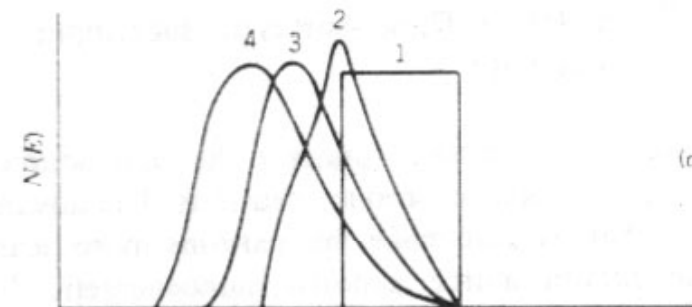
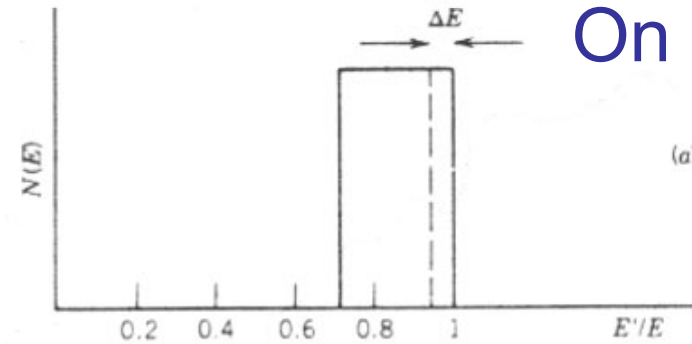
Average decrease in  $\ln(E)$  after one collision.

# Neutron Moderation

**HW 6 (continued)**

- Reproduce the plot. 
- Discuss the effect of the thermal motion of the moderator atoms.

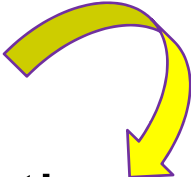
On  $^{12}\text{C}$ .



# Neutron Moderation

## HW 6 (continued)

Neutron scattering by light nuclei  $\overline{E'} = \frac{1}{2}(1 + \alpha)E$   
then the average energy loss  $\overline{\Delta E} = E - \overline{E'} = \frac{1}{2}(1 - \alpha)E$   
and the average fractional energy loss

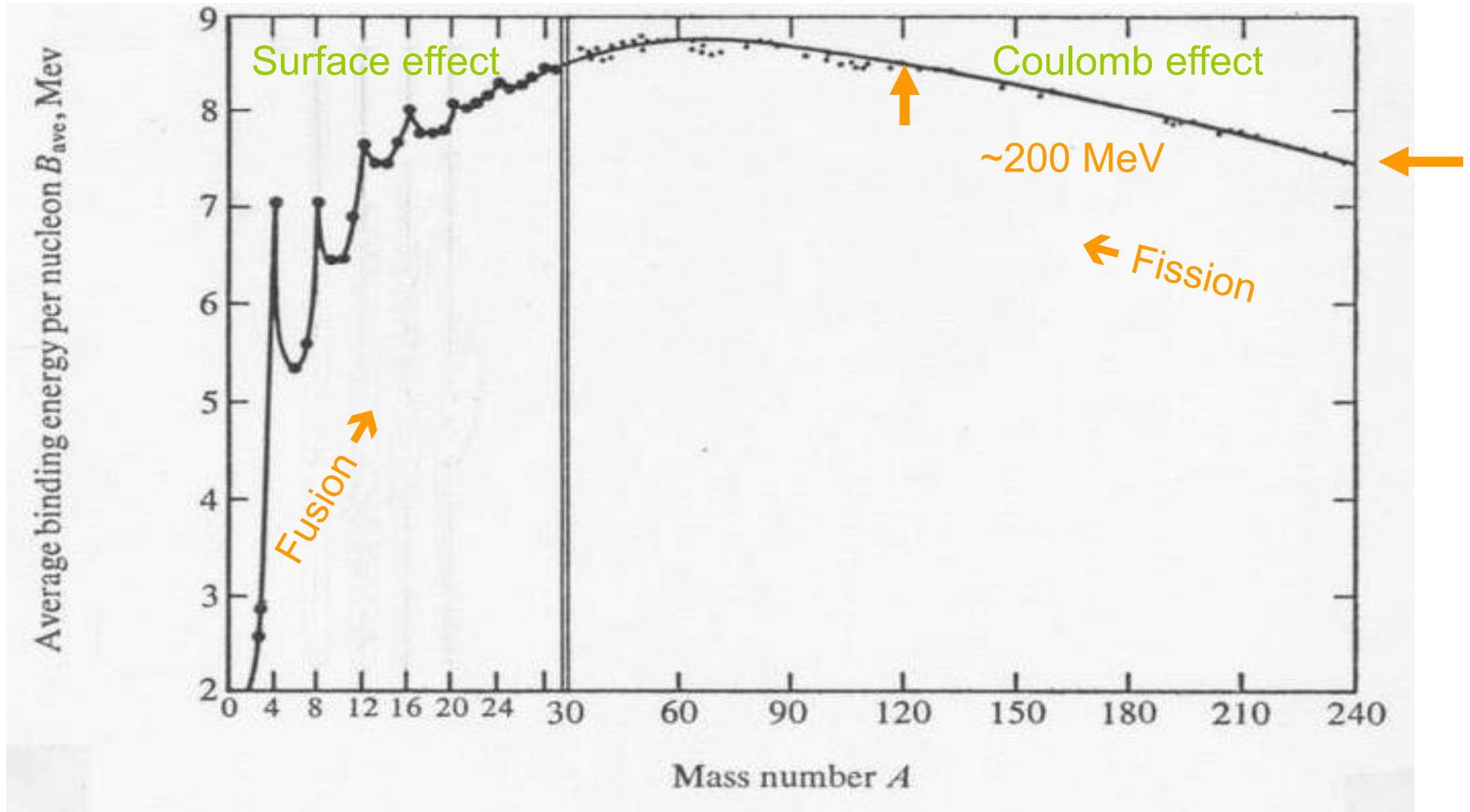
$$\frac{\overline{\Delta E}}{E} = \frac{1}{2}(1 - \alpha)$$


- How many collisions are needed to thermalize a 2 MeV neutron if the moderator was:

$^1\text{H}$     $^2\text{H}$     $^4\text{He}$    graphite    $^{238}\text{U}$    ?

- What is special about  $^1\text{H}$ ?
- Why we considered elastic scattering?
- When does inelastic scattering become important?

# Nuclear Fission



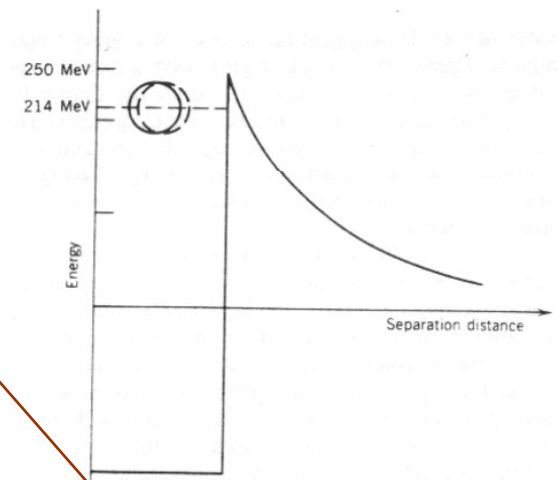
# Nuclear Fission

- B.E. per nucleon for  $^{238}\text{U}$  ( $\text{BE}_\text{U}$ ) and  $^{119}\text{Pd}$  ( $\text{BE}_\text{Pd}$ ) ?
- $2 \times 119 \times \text{BE}_\text{Pd} - 238 \times \text{BE}_\text{U} = ??$  ▶ K.E. of the fragments ▶  $\approx 10^{11} \text{ J/g}$
- Burning coal ▶  $10^5 \text{ J/g}$
- Why not spontaneous?
- Two  $^{119}\text{Pd}$  fragments just touching
- ▶ The Coulomb “barrier” is:

$$V = 1.44 \text{ MeV} \cdot \text{fm} \frac{(46)^2}{12.2 \text{ fm}} \approx 250 \text{ MeV} > 214 \text{ MeV}$$

- Crude ...! What if  $^{79}\text{Zn}$  and  $^{159}\text{Sm}$ ? Large neutron excess, released neutrons, sharp potential edge, spherical  $U$ ...!

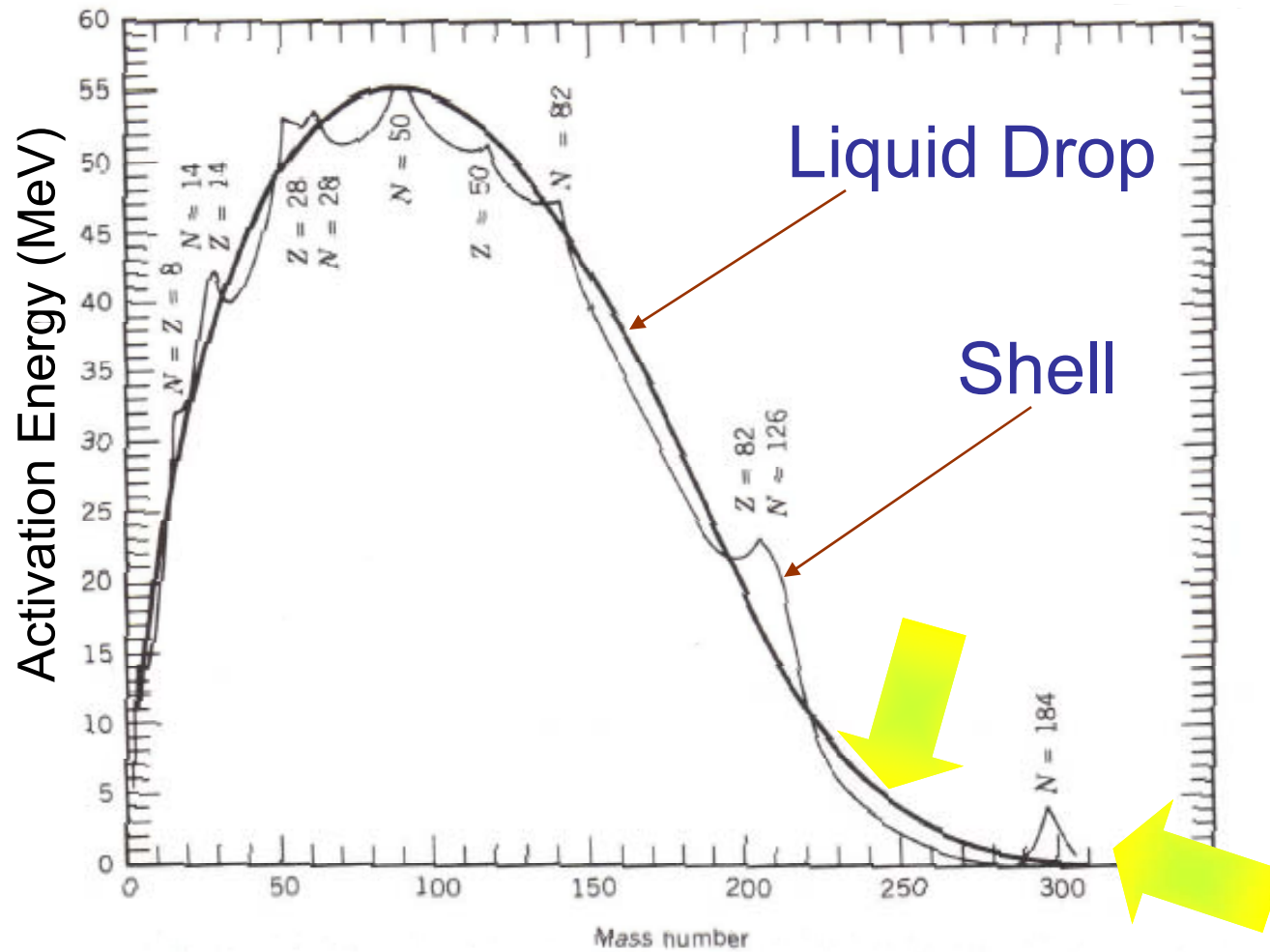
Crude



# Nuclear Fission

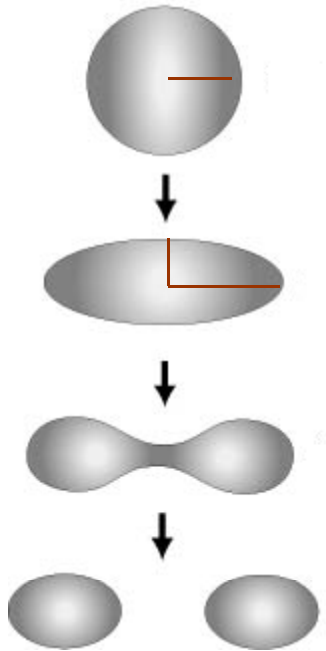
- $^{238}\text{U}$  ( $t_{1/2} = 4.5 \times 10^9$  y) for  $\alpha$ -decay.
- $^{238}\text{U}$  ( $t_{1/2} \approx 10^{16}$  y) for spontaneous fission.
- **Heavier nuclei??**
- Energy absorption from a neutron (for example) could form an intermediate state  $\blacktriangleright$  probably above barrier  $\blacktriangleright$  induced fission.
- Height of barrier is called activation energy.

# Nuclear Fission





# Nuclear Fission



$$\frac{4}{3}\pi R^3 = \frac{4}{3}\pi ab^2$$

$$R^3 = ab^2$$

$$a = R(1 + \varepsilon)$$

$$b = \frac{R}{\sqrt{1 + \varepsilon}}$$

Volume Term (the same)

Surface Term  $B_s = -a_s A^{2/3} (1 + \frac{2}{5}\varepsilon^2 + \dots)$

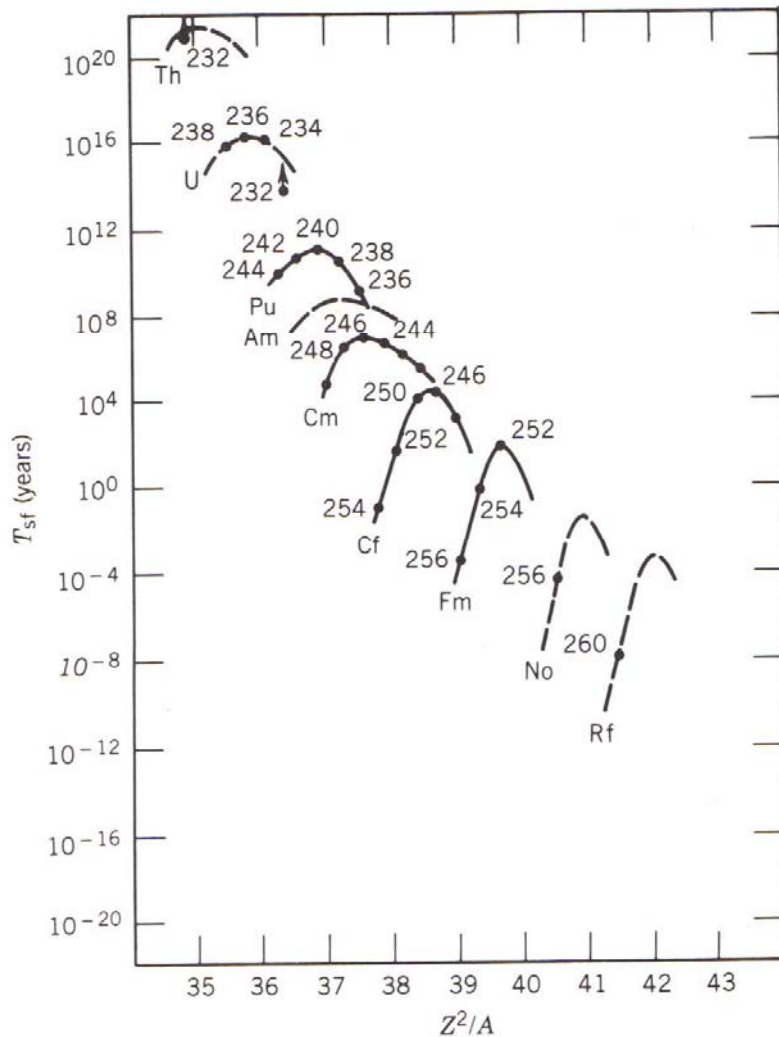
Coulomb Term  $B_C = -a_C Z(Z-1) / A^{1/3} (1 - \frac{1}{5}\varepsilon^2 + \dots)$

$$\frac{1}{5} a_C Z(Z-1) A^{-1/3} > \frac{2}{5} a_S A^{2/3} \blacktriangleright \text{fission}$$

$$\frac{Z^2}{A} > \sim 47$$

Crude: QM and original shape could be different from spherical.

# Nuclear Fission



$$\frac{(120)^2}{300} = 48$$

Consistent with activation energy curve for  $A = 300$ .

Extrapolation to 47  $\blacktriangleright \approx 10^{-20}$  s.