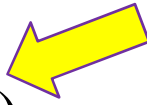


# Fuel Depletion

$$N \sim 10^{22} \text{ cm}^{-3}, \phi \sim 10^{14} \text{ cm}^{-2} \text{ s}^{-1}$$

**Time scale:  
Days and months.**

- More depletion ► increase steady state flux by means of reducing absorbers.

- For a given fuel isotope  $\frac{\partial N_f(\vec{r}, t)}{\partial t} = -N_f(\vec{r}, t)\sigma_a^f\phi(\vec{r}, t)$  

- For **constant flux**  $\phi_0$  the solution is **Exponential burnup**

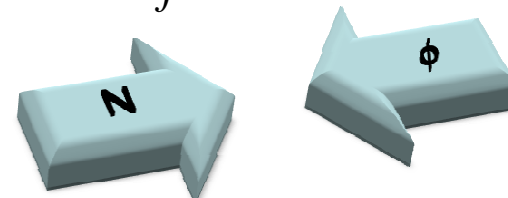
$$N_f(\vec{r}, t) = N_f(\vec{r}, 0)e^{-\sigma_a^f \phi_0(\vec{r})t} = N_f(\vec{r}, 0)e^{-\sigma_a^f \Phi(\vec{r}, t)}$$

- For **time varying flux**

$$N_f(\vec{r}, t) = N_f(\vec{r}, 0)e^{-\sigma_a^f \int_0^t \phi(\vec{r}, t') dt'} = N_f(\vec{r}, 0)e^{-\sigma_a^f \Phi(\vec{r}, t)}$$

**Neutron fluence**

**Solve numerically.**



# Fuel Depletion

- **Constant power.**

$$P(\vec{r}, t) = \overset{\downarrow}{w} N_f(\vec{r}, t) \overset{\uparrow}{\sigma_f^f} \phi(\vec{r}, t) = P(\vec{r}, 0) = P_0(\vec{r})$$

Energy released per fission  
Fission rate

$$N_f(\vec{r}, t)\phi(\vec{r}, t) = N_f(\vec{r}, 0)\phi(\vec{r}, 0)$$

$$\Sigma_f(\vec{r}, t)\phi(\vec{r}, t) = \Sigma_f(\vec{r}, 0)\phi(\vec{r}, 0)$$

- Power ~ flux only over short time periods during which  $N_f$  is constant.

$$\frac{\partial N_f(\vec{r}, t)}{\partial t} = -N_f(\vec{r}, t) \sigma_a^f \phi(\vec{r}, t) \approx -\frac{P_0(\vec{r})}{w}$$

Linear depletion!

- The solution is obviously

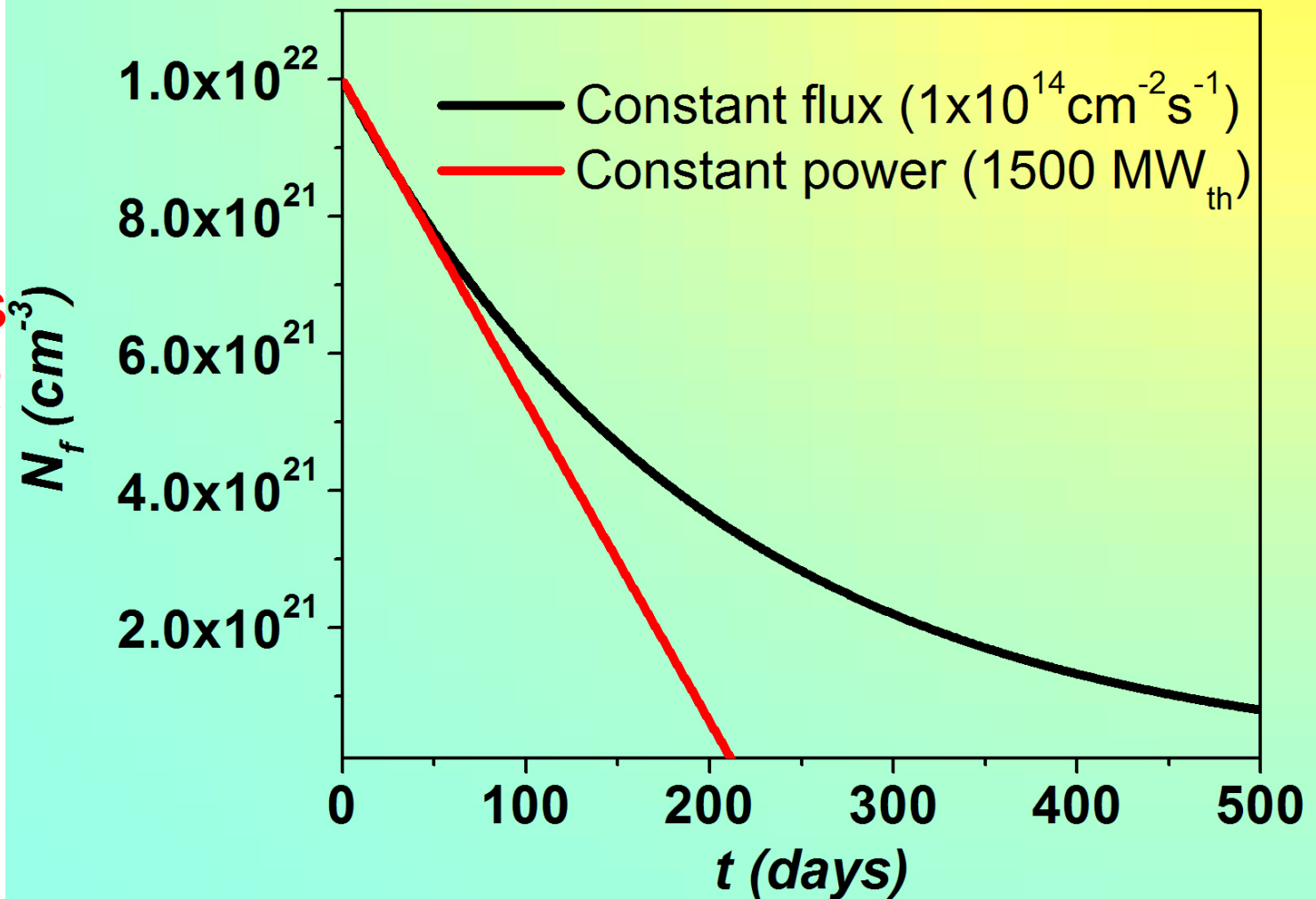
$$N_f(\vec{r}, t) \approx N_f(\vec{r}, 0) - \frac{P_0(\vec{r})}{w} t$$

$$\sigma_f^f \approx \sigma_a^f$$

# Fuel Depletion

## HW 31

Do the calculations for different flux and power levels.



# Poisoning and Fuel Depletion

Infinite, critical homogeneous reactor.

$$k_{\infty} = \epsilon \eta f \rho = \epsilon \eta \rho \frac{\sum_a^f(t)}{\sum_a^f(t) + \sum_a^{clad} + \sum_a^{moderator}(t) + \sum_a^{poison}(t) + \sum_a^{control}(t)}$$

$\downarrow$   $\downarrow$   $\uparrow$   $\downarrow$   
 thus  $\downarrow$

**Constant power**  $N_f(\vec{r}, t) \approx N_f(\vec{r}, 0) - \frac{P_0(\vec{r})}{w} t$

$$\phi(\vec{r}, t) = \frac{N_f(\vec{r}, 0)}{N_f(\vec{r}, t)} \phi(\vec{r}, 0) = \frac{\phi(\vec{r}, 0)}{1 - \sigma_a^f \phi(\vec{r}, 0) t}$$

$$= N_f(\vec{r}, 0) - N_f(\vec{r}, t) \sigma_a^f \phi(\vec{r}, t) t$$

$$= N_f(\vec{r}, 0) - N_f(\vec{r}, 0) \sigma_a^f \phi(\vec{r}, 0) t$$

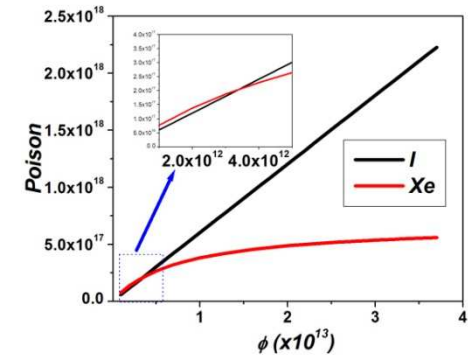
$$= N_f(\vec{r}, 0) [1 - \sigma_a^f \phi(\vec{r}, 0) t]$$

$$\sum_a^f(\vec{r}, t) = \sum_a^f(\vec{r}, 0) [1 - \sigma_a^f \phi(\vec{r}, 0) t]$$

# Poisoning and Fuel Depletion

$$Xe(t) = \frac{Xe_\infty}{\lambda_{Xe} + \sigma_a^{Xe} \phi_0} (1 - e^{-(\lambda_{Xe} + \sigma_a^{Xe} \phi_0)t}) + \frac{\gamma_I \sum_f \phi_0}{\lambda_{Xe} - \lambda_I + \sigma_a^{Xe} \phi_0} (e^{-(\lambda_{Xe} + \sigma_a^{Xe} \phi_0)t} - e^{-\lambda_I t})$$

Constant



$$\sum_a^{Xe}(\vec{r}, t) = \sigma_a^{Xe} Xe_\infty(\vec{r}, t) = \frac{(\gamma_I + \gamma_{Xe}) \sum_f(\vec{r}, 0) \phi(\vec{r}, 0)}{\frac{\lambda_{Xe}}{\sigma_a^{Xe}} + \phi(\vec{r}, t)}$$

Constant

$$\sum_a^{Sm}(\vec{r}, t) \approx \sigma_a^{Sm} \gamma_{Sm} \sum_f(\vec{r}, 0) \phi(\vec{r}, 0) t$$

$$\phi(\vec{r}, t) = \frac{\phi(\vec{r}, 0)}{1 - \sigma_a^f \phi(\vec{r}, 0) t}$$

- Other fission products (poisons) with less capture cross sections.

# Poisoning and Fuel Depletion

- Now we know all macroscopic cross sections.

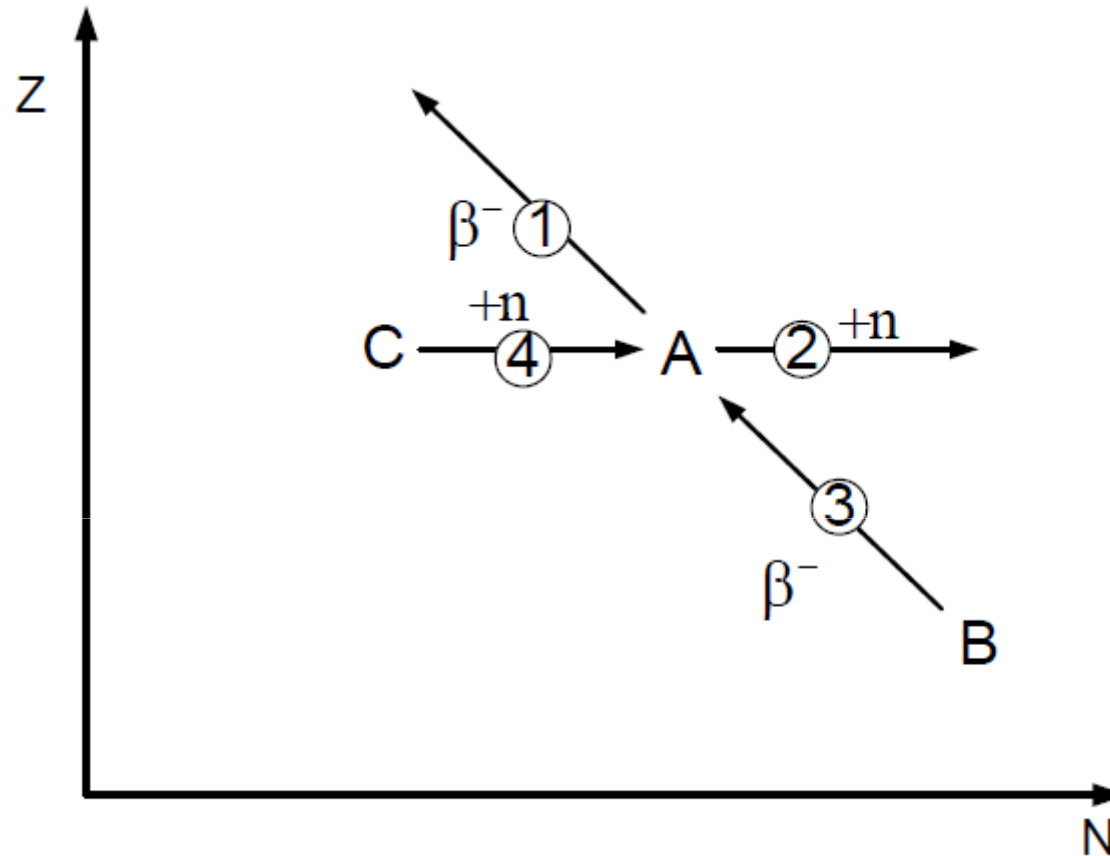
$$k_{\infty} = \varepsilon \eta f \rho = \varepsilon \eta \rho \frac{\sum_a^f(t)}{\sum_a^f(t) + \sum_a^{clad} + \sum_a^{moderator}(t) + \sum_a^{poison}(t) + \sum_a^{control}(t)}$$

↓
↑
↓

- When there are no absorbers left to remove, we need to refuel.
- Absorbers are not only control rods.
- All fuel nuclei should be considered.
- For each species, all sources and sinks should be taken into account.
- Online loading ► environmental.
- $^3\text{H}$ .

*Until = 0.*  
*Solve for t to get*  
*upper limit for*  
*“core loading*  
*lifetime”.*  
 Damaged  
 fuel...!

# Poisoning and Fuel Depletion



$$\frac{dN_A}{dt} = -\lambda_A N_A - \phi \sigma_A N_A + \lambda_B N_B + \phi \sigma_\gamma^C N_C + F(t)$$

Fuel loading

# Poisoning and Fuel Depletion

- **Some poisons are intentionally introduced into the reactor.**

- **Fixed burnable poisons.**

B, Gd.

More uniform distribution than rods, more intentionally localized than shim.

*Flatter flux.*

- **Soluble poisons (chemical shim) with caution.**

Boric acid (soluble boron, solbor) in coolant.

Boration and dilution.

Scram emergency shutdown (sodium polyborate or gadolinium nitrate).

- **Non-burnable poisons.**

Chain of absorbers or self shielding.

*Power shaping.*



# Delayed Precursors

$$\frac{1}{v_g} \frac{\partial}{\partial t} \phi_g(\vec{r}, t) = \chi_g \sum_{g'=1}^G \nu_{g'} \sum_{fg'}(\vec{r}) \phi_{g'}(\vec{r}, t) + \sum_{g'=1}^G \sum_{sg'g}(\vec{r}) \phi_{g'}(\vec{r}, t) + S_g^{ext}$$

$$- \sum_{ag}(\vec{r}) \phi_g(\vec{r}, t) - \sum_{sg}(\vec{r}) \phi_g(\vec{r}, t) + \vec{\nabla} \cdot D_g(\vec{r}) \vec{\nabla} \phi_g(\vec{r}, t)$$

- For one-group

$$\frac{1}{v} \frac{\partial}{\partial t} \phi(\vec{r}, t) = \nu \sum_f(\vec{r}) \phi(\vec{r}, t) + S^{ext}$$

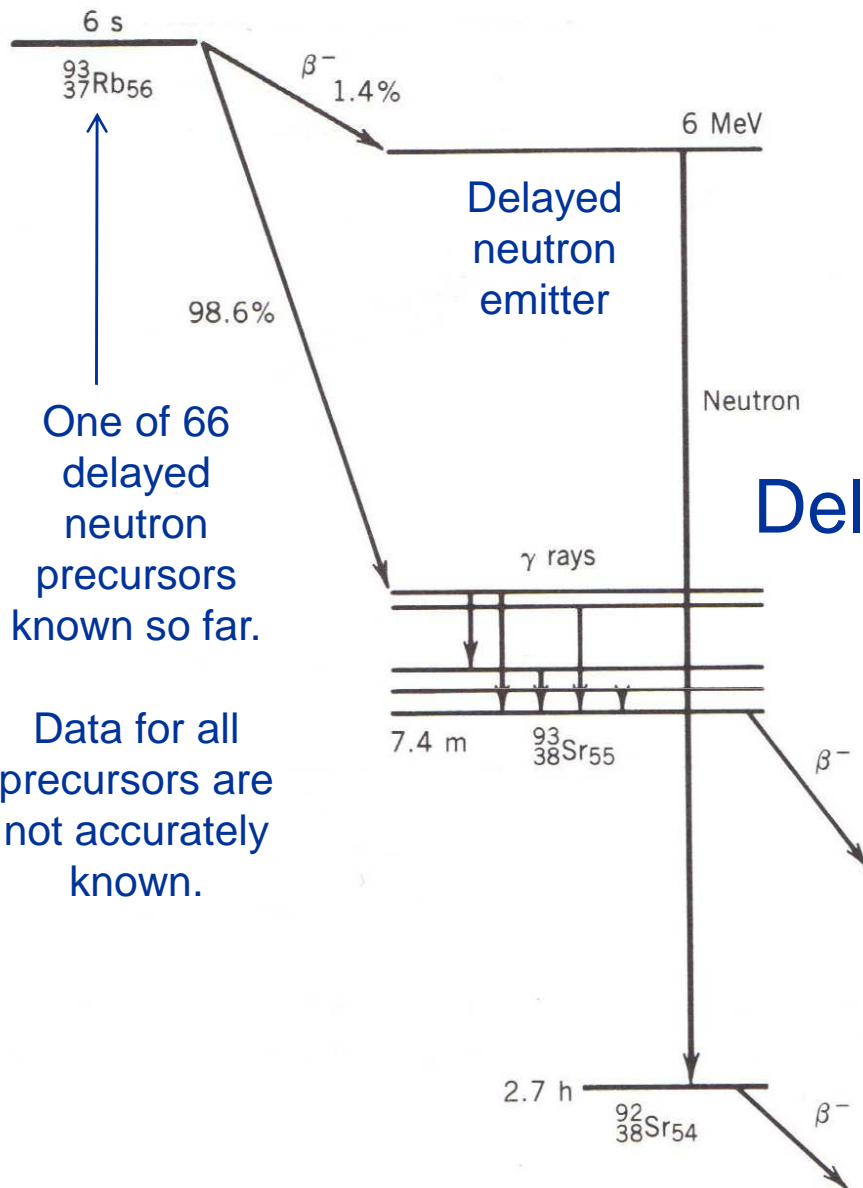
$$- \sum_a(\vec{r}) \phi(\vec{r}, t) + \vec{\nabla} \cdot D(\vec{r}) \vec{\nabla} \phi(\vec{r}, t)$$

- **What about delayed neutrons?**

# Delayed Precursors

$$\nu = \nu_p + \nu_d$$

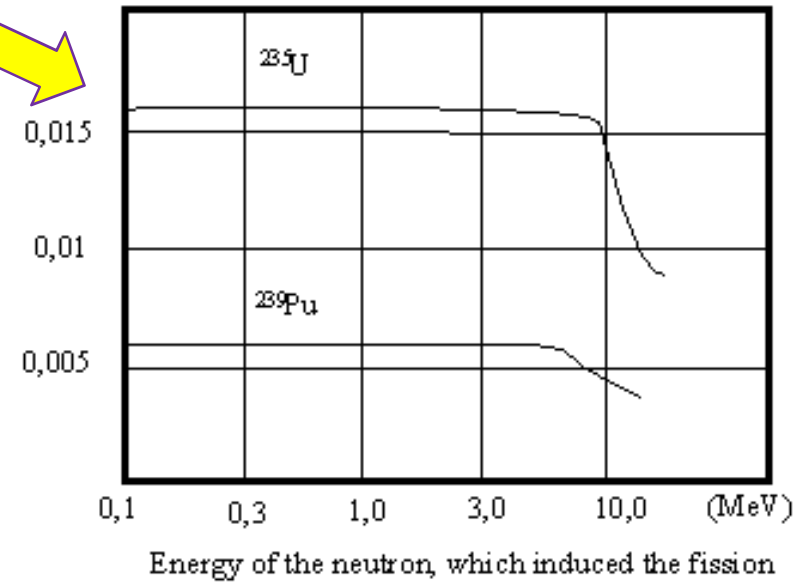
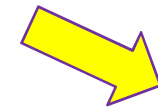
$$\text{Delayed neutron fraction } \beta = \frac{\nu_d}{\nu}$$



One of 66 delayed neutron precursors known so far.

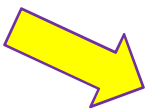

Data for all precursors are not accurately known.

Delayed neutron yield (delayed neutrons / fission)



# Delayed Precursors

| Fissile nucleus     | Delayed neutron / 100 fissions |
|---------------------|--------------------------------|
| $^{233}\text{U}$    | 0.667                          |
| $^{235}\text{U}$    | 1.621                          |
| $^{238}\text{U}^*$  | 4.39                           |
| $^{239}\text{Pu}$   | 0.628                          |
| $^{240}\text{Pu}^*$ | 0.95                           |
| $^{241}\text{Pu}$   | 1.52                           |
| $^{242}\text{Pu}^*$ | 2.21                           |

  Increases with  $N$ .

**Data for thermal neutron induced fission, except for  
\* fast neutron induced fission.**

# Delayed Precursors

| Group #→        | 1         | 2          | 3         | 4          | 5         | 6         |
|-----------------|-----------|------------|-----------|------------|-----------|-----------|
| $T_{1/2}$ (s)   | 54.51     | 21.84      | 6.0       | 2.23       | 0.496     | 0.179     |
| $\lambda_I$     | 0.0127    | 0.031      | 0.1155    | 0.310      | 1.397     | 3.871     |
| $\beta_i/\beta$ | .038      | 0.213      | 0.188     | 0.407      | 0.128     | 0.026     |
| $\beta_I$       | 0.0002641 | 0.00148035 | 0.0013066 | 0.00282865 | 0.0008896 | 0.0001807 |

$^{235}\text{U}$

$\nu_p$

$$\beta < 0.7\% = 0.016 / \nu$$

$$\frac{1}{\nu} \frac{\partial}{\partial t} \phi(\vec{r}, t) = (1 - \beta) \nu \sum_f(\vec{r}) \phi(\vec{r}, t) + \sum_{i=1}^6 \lambda_i C_i + S^{ext} - \sum_a(\vec{r}) \phi(\vec{r}, t) + \vec{\nabla} \cdot D(\vec{r}) \vec{\nabla} \phi(\vec{r}, t)$$

$$\frac{\partial C_i(\vec{r}, t)}{\partial t} = -\lambda_i C_i(\vec{r}, t) + \beta_i \nu \sum_f(\vec{r}) \phi(\vec{r}, t)$$

# Delayed Precursors

- The multi-group equation now becomes

Different energy spectra



$$\frac{1}{v_g} \frac{\partial}{\partial t} \phi_g(\vec{r}, t) = \chi_g^p (1 - \beta) \sum_{g'=1}^G v_{g'} \sum_{fg'}(\vec{r}) \phi_{g'}(\vec{r}, t) + \chi_g^c \sum_{i=1}^6 \lambda_i C_i(\vec{r}, t)$$

$$+ \sum_{g'=1}^G \sum_{sg'}(\vec{r}) \phi_{g'}(\vec{r}, t) + S_g^{ext}$$

$$- \sum_{ag}(\vec{r}) \phi_g(\vec{r}, t) - \sum_{sg}(\vec{r}) \phi_g(\vec{r}, t) + \vec{\nabla} \cdot D_g(\vec{r}) \vec{\nabla} \phi_g(\vec{r}, t)$$

$$\frac{\partial C_i(\vec{r}, t)}{\partial t} = -\lambda_i C_i(\vec{r}, t) + \beta_i \sum_{g'=1}^G v_{g'} \sum_{fg'}(\vec{r}) \phi_{g'}(\vec{r}, t)$$

# Full Blown Diffusion Equation

- In steady state

$$\lambda_i C_i(\vec{r}, t) = \beta_i \sum_{g'=1}^G \nu_{g'} \sum_{fg'}(\vec{r}) \phi_{g'}(\vec{r}, t)$$

$$0 = \chi_g^p (1 - \beta) \sum_{g'=1}^G \nu_{g'} \sum_{fg'}(\vec{r}) \phi_{g'}(\vec{r}) + \chi_g^c \beta \sum_{g'=1}^G \nu_{g'} \sum_{fg'}(\vec{r}) \phi_{g'}(\vec{r})$$

$$+ \sum_{g'=1}^G \sum_{sg'}(\vec{r}) \phi_{g'}(\vec{r}) + S_g^{ext} - \sum_{ag}(\vec{r}) \phi_g(\vec{r}) - \sum_{sg}(\vec{r}) \phi_g(\vec{r}) + \vec{\nabla} \cdot D_g(\vec{r}) \vec{\nabla} \phi_g(\vec{r})$$

Significance of  $\chi_g^c$  depends on whether we have fine or course energy groups.

$$0 = \left[ \chi_g^p + (\chi_g^c - \chi_g^p) \beta \right] \sum_{g'=1}^G \nu_{g'} \sum_{fg'}(\vec{r}) \phi_{g'}(\vec{r}) + \sum_{g'=1}^G \sum_{sg'}(\vec{r}) \phi_{g'}(\vec{r}) + S_g^{ext} - \sum_{ag}(\vec{r}) \phi_g(\vec{r}) - \sum_{sg}(\vec{r}) \phi_g(\vec{r}) + \vec{\nabla} \cdot D_g(\vec{r}) \vec{\nabla} \phi_g(\vec{r})$$